Massive MIMO Propagation Modeling With User-Induced Coupling Effects Using Ray-Tracing and FDTD

Sergei Shikhantsov, Arno Thielens, *Member, IEEE*, Günter Vermeeren, *Member, IEEE*, Piet Demeester, *Fellow, IEEE*, Luc Martens, *Member, IEEE*, Guy Torfs, *Senior Member, IEEE*, and Wout Joseph, *Senior Member, IEEE*

Abstract-The article describes a numerical approach for Massive MIMO channel modeling that accounts for the effects of electromagnetic coupling between a user and the receiving device. The modeling is performed by a combination of the Finite-Difference Time-Domain and the Ray-Tracing methods, supplemented with a stochastic geometry model of the propagation environment. The influence of user-coupling on the channel properties was studied statistically using the singular value spread and matrix power ratio metrics of the channel correlation matrix. The time-averaged Poynting vector distribution in the near-field of the receiver was evaluated using a realistic human phantom model and the Maximum Ratio Transmission precoding scheme in the downlink. The average enhancement of the timeaverage Poynting vector magnitude at the receiver location, compared to the surrounding area, was found to be around 10 dB when using 36 antenna elements at the base station. The electromagnetic field exposure of the phantom was assessed in terms of the 10g-average peak-spatial Specific Absorption Rate and compared with the existing public guidelines. Comparison of the EMF and exposure results provides a new perspective on the future regulatory procedures.

Index Terms—5G wireless, radio access networks, massive MIMO, propagation modeling, Ray-Tracing, FDTD, EMFexposure.

I. INTRODUCTION

RECENT years brought many advances to the development of next generation wireless networking. Significantly increased throughput, capacity and connection density are some of the requirements that 5G wireless technology must meet. Massive Multiple-Input Multiple-Output (MIMO) is one of the promising candidates to fulfill these requirements. Its operation is based on simultaneous utilization of a very large number of base-station (BS) antennas, (linear) transmission precoding and reception decoding schemes. An

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The authors are with the Department of Information Technology (INTEC), Ghent University–IMEC, 9052 Ghent, Belgium.

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accurate channel model is needed to facilitate its design and deployment. Both deterministic [1] and probabilistic [2], [3] approaches to Massive MIMO channel modeling are currently being developed.

One of the key features of realistic Massive MIMO channels is the scattering cluster spatial non-stationarity observed at the BS side [4], i.e. channel variability across the BS antenna elements. Another important effect is the variability of the receiver's radiation pattern due to its orientation in space and the effects of the near-field coupling with the user body.

In [5], the cluster non-stationarity is modeled introducing cluster visibility regions. However, effects on the user side are left unaccounted.

This contribution, to the best of the authors' knowledge, for the first time describes the numerical approach that takes both aforementioned effects of cluster non-stationarity and user body coupling into account. These effects will be utilized by the Massive MIMO downlink precoding schemes to produce compact space regions of an elevated electromagnetic field (EMF) around the receiver antennas. Therefore, we apply the proposed approach to calculate the EMF distribution in proximity of a user on a small scale, and use it to estimate the human EMF-exposure. The same method can also be adapted for a 5G near-user antenna design.

II. METHODS

The proposed approach is based on a hybrid Ray-Tracing (RT) method which is used to model large-scale propagation and the Finite-Difference Time-Domain (FDTD) method, used for refinement of the RT results. In the following sections each step of the approach is discussed in detail and the connection between them is explained.

A. Ray-Tracing

1) Method Description: The wireless channel modeling was performed using a ray-launching [6] variant of the RT method [7]. A commercially available REMCOM Wireless InSite 3.2 software package was used [8]. A general RT procedure relies on a ray-optics approximation of the Maxwell equations. A transmitter is modeled by launching rays from its center in a finite set of directions, distributed over the complete

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Fig. 1. Ray-Tracing environment model. The UE positions and the BS are shown; cuboid between them blocks direct propagation paths.

TABLE I Parameters of the RT Solver

Parameter	Reflections	Refractions	Diffractions	Threshold
Value	6	2	2	-80 dB(V/m)

unit sphere. A ray is propagated through the environment predefined with its geometry and material properties, undergoing reflections, refractions, and diffractions, until its power reaches a predefined threshold. If a ray passes in the vicinity of a receiver, it contributes to the total field at that receiver's location.

2) Environment Model: The RT environment model is shown in Fig. 1. It consists of a rectangular room, cuboid scatterers, a massive MIMO BS and receiver locations positioned on a straight line. The room has dimensions of 40 m×20 m×5 m. Dielectric properties $\varepsilon_r = 7$, $\sigma =$ $1.5 \cdot 10^{-2}$ S/m were assigned to its floor, ceiling and walls to model concrete material. The BS was modeled with a 6-by-6 planar rectangular array of uniform 1-wavelength (λ) interelement spacing. The BS elements were excited individually with a continuous sinusoidal signal at $f_c = 3.5$ GHz frequency ($\lambda \simeq 86$ mm). The array plane was normal to the x-axis and its center was set at at x = 7 m, y = 10 m, z = 4 m, as indicated in Fig. 1.

The 19 UE locations were arranged in a straight line along the x-axis separated by 1 m, in a range from 15 m to 33 m at the height of 1 m.

A perfect electrical conductor (PEC) cuboid of size $2 \text{ m} \times 0.2 \text{ m} \times 4 \text{ m}$ (around $20\lambda \times 2\lambda \times 40\lambda$ at 3.5 GHz) was placed between the BS and the UEs (x = 10 m, y = 10 m), thus blocking the line-of-sight (LOS) propagation. Other PEC cuboid scatterers ($2 \text{ m} \times 0.5 \text{ m}$ size and height from 2 m to 3 m) were randomly distributed along the perimeter of the room and had uniform random rotation along the vertical axis. This model generates realistic industrial environments of a fixed layout (e.g. a warehouse, assembly line) [9].

A 0.02° ray-spacing was set for the ray-launching at the Tx. Each antenna element had vertical polarization and used an isotropic radiation pattern. Other RT solver parameters are summarized in Table I and were chosen as suggested in [10].

A single realization of the random scatterers produces an *environment sample*. By successively generating environment samples with random arrangement of scatterers and tracing

rays in these samples, we obtained statistics of the EM-field that was incident at the NLOS locations in the described environment. The EM-field properties crucial for the operation of the massive MIMO (e.g. correlation between signals emitted from different antennas) are calculated based on a realistic geometry.

For each Tx-Rx pair (n, k) in the simulation, the RT solver returned a collection of rays $\{\mathbf{r}\}_{n,k}$. A ray holds properties of the propagation path it models, such as direction-of-arrival (DOA), time-of-flight (TOF), Path Loss (PL), EM-field strength, etc. We used these properties to calculate the wireless channel in Section II-C.

B. EMF Coupling Effects

This section explains how the radiation pattern of the UE close to a phantom's head can be calculated using FDTD simulations.

In realistic scenarios, such as a mobile phone call, a user penetrates the near-field of his UE. This results in EMcoupling which affects the radiation pattern of the UE. We accounted for this by simulating the UE antenna in a usage mode together with a realistic human body model using FDTD method.

To model the UE antenna, we used a generic halfwavelength dipole antenna model, designed for the central frequency of 3.5 GHz. The overall dipole antenna length was 37 mm, arm diameter and the feed gap were 2 mm. The dipole arms were modeled as perfect electric conductors.

We used the ViP v.3.1 heterogeneous Duke human phantom [11] (further referred as *the phantom*) as a human model.

The Sim4Life v4.4 (Zürich, Switzerland) software package was used for FDTD simulations throughout the paper. Fig. 2 shows the FDTD computational domain. The antenna was positioned at the center of the global coordinate system and oriented vertically. The phantom's head was positioned near the dipole, such that its bounding box center was coincident with the y-axis and the distance from the left ear to the dipole center was 20 mm. This setup aimed to reproduce the effects observed in a real-life phone call scenario, such as signal blockage by the user head.

The dipole was fed with a sinusoidal signal at 3.5 GHz and 1 W of total input power. The output of the simulation was a complex amplitude of the vertically polarized far-field E-field $\hat{A}(\theta, \phi)$. $\hat{A}(\theta, \phi)$ was sampled on a surface of a 1 m radius sphere for elevation angle θ in range $[0, \pi]$ and azimuth angle φ in range $[0, 2\pi)$ with a 2° step. Its ratio to the far-field E-field strength of an isotropic radiator fed the same input power (\simeq 7.74 V/m) is the normalized far-field $A(\theta, \phi)$, which will be used throughout the following sections.

The computational domain center was coincident with the center of the phantom's head bounding box and its dimensions were set to $256 \text{ mm} \times 256 \text{ mm} \times 250 \text{ mm}$. It fully enclosed the head and the antenna, at the same time significantly reducing the computational demands needed for simulations described in the next section. A maximum discrepancy of around 8% was observed in the antenna directivity if compared to the full-body simulation, which was considered acceptable, taking into account a simplified model of the UE.



Fig. 2. The FDTD simulation setup. Solid black lines show the domain boundaries. The original model of ViP v.3.1 Duke phantom's head is presented alongside with its voxelized approximation. The dipole near the phantom's left ear is shown in blue together with its radiation pattern. Global Cartesian and spherical coordinate systems used throughout the paper are given.

C. Channel Matrix

This section explains how the channel matrix was calculated using rays at the Rx and the radiation pattern from the FDTD simulations.

First, a free-space channel matrix \mathbf{H}^{fs} was constructed from the collection of all rays $\{\mathbf{r}\} = \bigcup \{\mathbf{r}\}_{n,k}$. A channel between the Tx antenna with index n and the Rx antenna with index k is found as

$$h_{k,n}^{fs} = \sum_{r \in s(n,k)} \tilde{E}_r^{\theta}, \tag{1}$$

and

$$\tilde{E}_r^{\theta} = E_r^{\theta} \exp(-2\pi i f_c \tau_r).$$
⁽²⁾

Here s(n,k) are the indices of the rays in $\{\mathbf{r}\}_{n,k}$, E_r^{θ} and τ_r are the vertical polarization component of the E-field and the TOF of the r^{th} ray respectively.

Second, a reciprocity of an antenna in transmit-receive was utilized to introduce the DOA dependence into (1). As a ray holds the information about its DOA (θ_r, φ_r), by weighting its contribution to the channel in (1) with the corresponding radiation pattern value $A(\theta_r, \varphi_r)$, coupling effects are introduced into the channel matrix elements

$$h_{k,n}^{nf} = \sum_{r \in s(n,k)} A(\theta_r, \varphi_r) \tilde{E}_r^{\theta}.$$
(3)

For an arbitrary incident direction (θ_r, φ_r) , $A(\theta_r, \varphi_r)$ is calculated by the bilinear interpolation of $A(\theta, \varphi)$. Evaluating (3) for every Tx-Rx pair, we obtain a full massive MIMO channel matrix \mathbf{H}^{nf} .

We calculated the Singular Value Spread (SVS) κ (G) and Matrix Power Ratio (MPR) γ (G) of the channel correlation matrix G to quantify how the inclusion of the UE radiation pattern into the model affects the channel. The channel correlation matrix is obtained as

$$\mathbf{G} = \mathbf{H}\mathbf{H}^{H},\tag{4}$$

where $(\cdot)^H$ denotes the conjugate transpose of a matrix.

The SVS of the channel correlation matrix is widely used in MIMO analysis [12]; it is a ratio between the largest and the smallest singular values of **G**.

The MPR is the ratio between the sum of the squared magnitudes of diagonal elements of G and all its elements [13]

$$\gamma(\mathbf{G}) = \frac{\sum_{i=1}^{K} |g_{i,i}|^2}{\sum_{i=1}^{K} \sum_{j=1}^{K} |g_{i,j}|^2},$$
(5)

where $g_{i,j}$ are elements of **G** and *K* is the number of receivers.

D. EMF Distribution

To determine the EMF distribution in proximity of the user and the UE, FDTD simulation based on the RT results were performed. A ray with index r was modelled as a planewave source $\mathbf{p}_r = (\mathbf{k}_r, a_r, \phi_r)$ that spans across the complete computational domain. It is described by its wave-vector \mathbf{k}_r , amplitude a_r and phase delay at the domain center ϕ_r .

To reduce the required computational resources, DOAs of the rays in **r** were substituted by the outer normal vectors of the faces of an icosahedral sphere (ico-sphere) $\{\mathbf{n}_{ico}(m)\}$ [14], with *m* being the ico-sphere frequency. The wave-vector of a plane wave that corresponds to the *i*th face of the ico-sphere is then the inner normal of that face

$$\mathbf{k}_i = -\mathbf{n}_{ico,i}.\tag{6}$$

The amplitude and phase of the plane wave with index i were obtained by taking a sum of the complex amplitudes of the rays, DOAs of which have the smallest angular distance to the i^{th} ico-sphere outer normal. Each ray's amplitude was weighted with the *precoding matrix* element $w_{n,k}$, where (n, k) is the Tx-Rx pair indices for which the ray was calculated.

In this contribution we investigated the Maximum Ratio Transmission (MRT) precoding scheme [12]. MRT matrix W is proportional to the complex conjugate transpose of the channel matrix and its elements are given by

$$w_{n,k} = \alpha h_{k,n}^*,\tag{7}$$

where the normalization coefficient α is chosen such that W has unit Frobenius norm.

The amplitude and phase of the plane wave with index i at the Rx with index k is then given by the magnitude and argument respectively of

$$p_{i,k} = \sum_{n=1}^{N} w_{n,k} \sum_{r \in s(n,k)} \tilde{E}_{r}^{\theta}.$$
 (8)

Here the outer sum is taken over the BS elements, with N being their overall number.

Using (7) to calculate the precoding with (3) as the channel matrix, we obtained a set of the plane wave sources that models the EMF incidence at a massive MIMO user

$$p_{i,k}^{nf} = \alpha \sum_{n=1}^{N} \sum_{r \in s(n,k)} (h_{k,n}^{nf})^* \tilde{E}_r^{\theta}.$$
 (9)

One FDTD simulation was required for each user in the environment. The result of a simulation was a 3-dimensional distribution of the electromagnetic field in free-space around the phantom and inside its tissues. Knowing the dielectric properties of the phantom's tissues we were able to calculate its EMF-exposure in terms of the Specific Absorption Rate (SAR). We used the peak-spatial SAR averaged over a 10-g cube (psSAR_{10g}), calculated according to the IEEE/IEC 62704-1 standard [15]. psSAR_{10g} captures local exposure peaks, that highly-focused fields distinctive to the massive MIMO technology, are expected to produce.

International Comission on Non-Ionizing Radiation (ICNIRP) specifies basic restrictions for $psSAR_{10g}$ [16] in the head for the general public (2 W/kg) and workers (10 W/kg). Based on these limits, reference levels on time-averaged power density are established using FDTD simulations with a single incident plane wave (10 W/m² and 50 W/m² for the general public and workers respectively). We will compare massive MIMO exposure and free-space power density with the ICNIRP guidelines in the following section.

For a more detailed description of the FDTD simulation setup, its sensitivity and error analysis, see [9].

III. RESULTS

In this section, the results are given in the order of the subsections of the previous section.

A. Far-Field Pattern

The normalized radiation pattern $A(\theta, \varphi)$ was calculated as described in Section II-B. Fig. 3 shows its magnitude and phase in spherical coordinates. Azimuth angle $\varphi = \pi/2$ corresponded to the incident radiation that is not obstructed by the user's head (see the coordinate system at Fig. 2).

A global maximum of |A| was observed around this azimuth angle in the horizontal plane ($\theta = \pi/2$). This was expected, as the UE dipole is vertically oriented, it favours propagation in the horizontal plane. The phase response shown at the bottom of Fig. 3 is relatively flat around the magnitude maximum, varying for no more than $\pi/4$ in the $\pi/2$ neighbourhood of ($\theta = \pi/2, \varphi = \pi/2$).

A global minimum of A magnitude was found near $\phi = 3\pi/2$. This was also expected, as at these angles the incident radiation is attenuated by the user's head. The amount of attenuation is significant; it reaches around -18 dB (a factor of 63) if compared to the global maximum.

The phase of A oscillated rapidly around $\phi = 3\pi/2$. This can be explained by a superposition of multiple diffraction



Fig. 3. Radiation pattern of the dipole antenna coupled with the phantom's head. Top: Amplitude. Bottom: Phase.

paths that become dominant when the LOS is blocked by the head. However, this rapid phase variation does not make any noticeable contribution to the channel due to a low relative power of the propagation paths associated with it.

B. Channel Correlation Matrix

100 environment samples were simulated using the RT as described in Section II-A2.

First, a free-space channel correlation matrix \mathbf{G}^{fs} was calculated. This was done by evaluating (4) with \mathbf{H} given by (1) for rays traced in each environment sample. At the top of Fig. 4 the magnitude of a sample average of \mathbf{G}^{fs} is shown.

Second, the channel correlation matrix \mathbf{G}^{nf} , that accounts for the UE radiation pattern was calculated. For each UE in the environment, the radiation pattern $A(\theta, \varphi)$ (see Section II-B) was rotated around the z-axis to an angle sampled randomly in $[0, 2\pi)$. This models a random positioning of users with respect to the BS, when all user orientation directions were equally probable. The average \mathbf{G}^{nf} taken over all 100 environment samples is shown at the bottom of Fig. 4.

The diagonal elements of **G** were proportional to the power received in the downlink (DL) by users, if the MRT precoding was applied. Both \mathbf{G}^{fs} and \mathbf{G}^{nf} were diagonally dominant. This means that a significant portion of the DL transmitted power reached the intended receiver. However, in case of \mathbf{G}^{fs} there was an apparent increase in a relative magnitude of the super- and sub-diagonal elements. This corresponds to the increased interference of users with their closest neighbours in the simulated environment. Such effect was not observed in \mathbf{G}^{nf} . An average ratio between the elements is around 12% for \mathbf{G}^{fs} and less that 8% for \mathbf{G}^{nf} .

This difference (around 30% decrease) can be explained by an elevated correlation in the incident field as a function of DOA between closely separated locations. Presence of a highly-directional individual radiation patterns helps to



Fig. 4. Normalized channel correlation matrix averaged over 100 environment samples (logarithmic scale). Top: Free space channel \mathbf{H}^{fs} . Bottom: Channel with the coupling effects \mathbf{H}^{nf} .

resolve users that do not have enough spatial separation, which will have a positive effect on a massive MIMO system performance.

To study statistical properties of the channels, we calculated MPR and SVS for a different numbers of active UEs in the environment and Tx elements used at the BS. First, two and five UEs were selected randomly from the full simulated set and channel matrices (1) and (3) were calculated with only those UEs (index k) and a full 36-element array at the BS. Second, the same calculation was performed with five and ten randomly selected BS antenna elements and five randomly selected UEs. This was repeated 100 times in 100 environment samples, resulting in 10^4 samples per dataset. Additionally, $\kappa(\mathbf{G})$ and $\gamma(\mathbf{G})$ of the full channel (19 UEs) were calculated.

Fig. 5(a) depicts histograms of $\kappa(\mathbf{G})$ and $\gamma(\mathbf{G})$ for 2, 5 and 19 UEs in black, blue and red, respectively. As the number of users in the channel increases, MPR (left column) decreases for both \mathbf{G}^{fs} (top) and \mathbf{G}^{nf} (bottom). This is in line with definition (5), as the number of the off-diagonal elements of a square matrix is roughly proportional to the square of the number of its diagonal elements. However, $\gamma(\mathbf{G}^{nf})$ has a



(a) Distributions for 2 and 5 randomly selected UEs, and all 19 UEs are shown with black, blue, and red respectively. All 36 Tx elements are used.



(b) Distributions for 4 and 9 randomly selected Tx elements, and all 36 Tx elements are shown with black, red, and blue respectively. 5 UEs are randomly selected in each channel evaluation.

Fig. 5. Matrix Power Ratio $\gamma(\mathbf{G})$ (left column) and Singular Value Spread $\kappa(\mathbf{H})$ (right column) empirical distributions of 100 environment sample channels. In each sub-figure free-space channels (top rows) and channels with the coupling effects (bottom rows) are shown.

larger positive skew and average values compared to $\gamma(\mathbf{G}^{fs})$. Higher $\gamma(\mathbf{G})$ corresponds to channels with less correlation between UEs, being unity for purely orthogonal channels (with the correlation matrix \mathbf{G} being diagonal). This means that the near-field user-UE coupling decreases (on average) correlation between massive MIMO users.

The right column in Fig. 5(a) shows $\kappa(\mathbf{G})$ (in dB) for \mathbf{G}^{fs} (top) and \mathbf{G}^{nf} (bottom). For both \mathbf{G}^{fs} (top) and \mathbf{G}^{nf} larger user count yields larger SVS values. This was expected, as $\kappa(\mathbf{G})$ is determined by only two most correlated rows of \mathbf{G} . The more UEs are included in the channel (the greater dim(\mathbf{G}) is), the higher chance there is for any two users to have correlated channels.

Fig. 5(b) presents histograms of $\kappa(\mathbf{G})$ and $\gamma(\mathbf{G})$ for a fixed number of 5 UEs and a varying number of the BS antenna elements. With an increase of the element count, the mean of $\gamma(\mathbf{G})$ monotonically grows for both \mathbf{G}^{fs} and \mathbf{G}^{nf} . The mean of $\gamma(\mathbf{G}^{nf})$ remains larger than $\gamma(\mathbf{G}^{fs})$ for all (equal) studies BS antenna counts.

Increasing the number of the BS antennas leads to a decrease of $\kappa(\mathbf{G})$, as shown in the right column of Fig. 5(b). Moreover, the amount of $\kappa(\mathbf{G})$ variation strongly depends on the ratio N/K between the number of the BS antennas and

TABLE II Mean Values of $\kappa(\mathbf{G})$ and $\gamma(\mathbf{G})$



Fig. 6. Magnitude of the time-averaged Poynting vector distribution. Slice in the horizontal plane (z = 0) shows the average of 19 user locations in 10 environment samples.

the UEs. Increasing N from 5 to 10 (M/K from 1 to 2) leads to around 6 dB (4 times) decrease of mean values of both $\kappa(\mathbf{G}^{nf})$ and $\kappa(\mathbf{G}^{fs})$. Increasing N from 10 to 36 (N/K from 2 to 6.2) decreases the mean value of $\kappa(\mathbf{G}^{fs})$ around 1 dB. This observation agrees well with the results of measurements in NLOS conditions [17].

Mean values of the histograms shown at Fig. 5 are compiled in Table II. In summary, the near-field coupling leads to the increase in the average of both $\gamma(\mathbf{G})$ and $\kappa \mathbf{G}$.

C. Hot-Spot

In this section we present the results of the FDTD simulations described in Section III-C.

We performed FDTD simulations for all UEs in 10 environment samples. To evaluate the focusing of the EMF in proximity of the UE antenna we calculated the time-averaged Poynting vector at the receiver k as

$$\mathbf{S}_k = \operatorname{Re}(\frac{\mathbf{E}_k \times \mathbf{H}_k^*}{2}),\tag{10}$$

where electric and magnetic field vectors \mathbf{E}_k and \mathbf{H}_k were interpolated on a rectilinear grid. A horizontal slice coincident with the phantom's head center (z = 0) of $||\mathbf{S}_k||$, averaged over all 19 UEs in 10 environment samples is shown at Fig. 6.

A strong focusing of $||\mathbf{S}_k||$ is present near the dipole center, shown with a black circle in Fig. 6. This was the compound effect of the MRT precoding and the propagation environment. The precoding forced the signals from different BS elements to arrive coherently at the antenna terminal and a sufficiently



Fig. 7. Sample density of left/right time-averaged Poynting vector gain.

diverse environment makes it very unlikely for them to be in-phase at any other point in space.

To quantify the focusing effect we calculated the ratio between time-averaged power density at the antenna center $\mathbf{r}_{Rx} = (0, 100, 0)$ (shown with a circle) and $\mathbf{r}_{sym} = (0, -100, 0)$ that was symmetric to it with respect to the *xz*-plane (marked with a black cross at Fig. 6)

$$S_k = \frac{\|\mathbf{S}_k(\mathbf{r}_{Rx})\|}{\|\mathbf{S}_k(\mathbf{r}_{sym})\|}.$$
(11)

As both the environment model and the phantom body were symmetric with respect to the xz-plane, it is reasonable to assume that for DL transmission with no precoding, the average of S_k would be 1.

Fig. 7 shows a histogram of S_k (dB scale). Average S_k is around 10 dB, and 5th - 95th percentile range spans from just above 3 to slightly over 16 dB. This means that that the average hot-spot values of the power flux density are 10 times higher than in the surrounding space.

D. Specific Absorption Rate

The psSAR_{10g} was calculated in the phantom's head using built-in algorithms available in the Sim4Life software. We evaluated SAR_{10g} in the same environment samples and UE locations as were analyzed in Section III-C, resulting in 190 overall number of exposure samples. Fig. 8 depicts SAR_{10g} distribution in the horizontal slice at z = 0 and projections of the peak-cubes on that plane.

The largest portion of the peak-cubes ($\simeq 78\%$) were found near the left ear (see Fig. 8, red circle). These exposure peaks are produced by the EMF hot-spot at the UE antenna, located close to the ear.

In contrast to that, only around 9% of all peak-cubes is found at the right ear (indicated with the green circle at Fig. 8). These peak-cubes are spatially more concentrated than the cubes at the left ear, which suggests that they were mainly caused by the ear geometry. High curvature of the ear tends to concentrate EMF in the surrounding tissues even if exposed to a plane-wave [18]. This effect can be observed in Fig. 6 by noting the in-tissue local maximas of the power flux density. The most of the remaining peak-cubes were located at the back of the head, in the direction to the BS. In three samples the psSAR_{10g} was found in the eyes.

This allows to conclude that most of the Massive MIMO DL exposure was produced in the vicinity of the UE.



Fig. 8. psSAR_{10g} and peak-cubes in the phantom's head. Slice in the horizontal plane (z = 0) shows the average of 19 user locations in 10 environment samples.

To perform the exposure analysis, we introduced η - the psSAR_{10g} normalized to the time-averaged Poynting vector magnitude. Using the normalized exposure aids the analysis in two aspects. First, the influence of the BS transmit power and the PL is eliminated. Thus, values assessed at the UEs located at different distances to the BS can be correctly compared. Second, η is an indicator of how adequate the existing safety regulatory limits are. Current regulations require the EMF measured in free-space to comply with a predefined threshold. The threshold value was established based on the results of extensive direct measurements, simulations and additional safety factors, in which conventional single-antenna transmitters were used.

The definition of η depends on the location and conditions at which the EMF is assessed for normalization. Two normalization strategies were compared. First, η_{hs} was calculated as psSAR_{10g} normalized to the maximum $\|\mathbf{S}_k\|$ in the domain (hot-spot normalization). This is the time-averaged power flux density, that would occur close to the UE in operation.

Second, η_{pw} was calculated by normalizing psSAR_{10g} to the average of the power flux densities in the incident planewave sources defined in (9). This is the value that would be observed in free-space when no user is present, as currently followed by the ICNIRP exposure guidelines [16].

As a reference we used $\eta_{ref} = 0.2 \text{ m}^2/\text{kg}$, calculated from the ICNIRP general public restrictions (psSAR_{10g}^{ICNIRP} = 2 W/kg, $\|\mathbf{S}\|^{ICNIRP} = 10 \text{ W/m}^2$).

Fig. 9 depicts mean values and 5th - 95th percentiles of η_{hs} and η_{pw} calculated for 190 exposure samples, plotted versus the distance to the BS. For all studied distances η_{hs} was nearly constant and around 3 times lower than η_{ref} . It also had a relatively low variance: 90% of all samples fell between -3 and -7 dB. Additionally, psSAR_{10g} was highly correlated with the hot-spot power flux density (Pearson correlation coefficient $\rho = 0.96$). This was in agreement with the analysis of the peak-cube locations: most of the exposure was observed near the hot-spot. Hence, power flux density close to the



Fig. 9. psSAR_{10g} normalized to the time-averaged Poynting vector (dB scale). η_{hs} (blue)-normalization using values at the hot-spot. η_{pw} (green)-normalization with the average of the incident plane waves. For all UE locations η_{hs} is lower and η_{pw} is higher than the value calculated from the ICNIRP basic restrictions (used for the dB reference).

receiver could be used to estimate the EMF-exposure. In addition, applying existing reference levels would overestimate $psSAR_{10g}$ at least by a factor of 2.

In contrast, the average η_{pw} is about 15 dB higher than η_{ref} and has larger sample variation compared to η_{hs} . This indicates that free space EMF measurements cannot reliably estimate the massive MIMO exposure. The existing free-space reference levels underestimate psSAR_{10g} by a factor ranging from 10 to 100.

IV. CONCLUSION AND FUTURE WORK

A novel numerical approach to massive MIMO channel modeling based on the RT and the FDTD methods was presented. For the first time, to the authors' knowledge, massive MIMO channels were simulated with the effects of near-field coupling between the receiver antenna and the user body. The importance of this effect was studied with models of a generic dipole antenna, a realistic human phantom and an industrial indoor environment. It was shown that including the coupling effects decreases correlation between closely spaced UEs by around 30%. The time-averaged Poynting vector magnitude enhancement at the UE was around 10 dB. psSAR_{10g} in the phantom's head was found to be directly proportional to the hot-spot power flux density. Normalized psSAR_{10g} complies with the ICNIRP guidelines, but assessing the reference levels in free space (with no active receiver) would lead to its underestimation by at least a factor of 10.

Although the presented study was carried out at a sub-6 GHz frequency, the same approach can be applied to a mm-Wave system analysis. The RT method is more accurate at higher frequencies with no performance loss. The FDTD memory grows as the third power, and time as the first power of frequency, which can limit the overall performance.

Future work will focus on studying the coupling between the BS antenna elements on the one hand and the user and the UE on the other hand, as these effects could have a significant impact on the channel [19].

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Sergei Shikhantsov received the B.Sc. and M.Sc. degrees in applied physics and mathematics from the Moscow Institute of Physics and Technology (MIPT), Moscow, Russia, in 2014 and 2016, respectively. He is currently pursuing the Ph.D. degree in engineering physics with Ghent University, Ghent, Belgium. His research interests include computational electrodynamics, numerical assessment of human electromagnetic field exposure, and propagation modeling of next generation wireless networks.



Arno Thielens (Member, IEEE) received the M.Sc. (Eng.) and the Ph.D. degrees in applied physics from Ghent University in September 2010 and May 2015, respectively. He joined the Waves group, at that time known as WiCa, of Ghent University's Information Technology (INTEC) Department, as a Ph.D. student in May 2015. Since June 2015, he has been working as a Post-Doctoral Researcher at Waves. He was a Post-Doctoral Fellow at the Institute for Science and Innovation (IWT), Flanders, Belgium, from January 2016 to May 2017. In June 2017, he joined

Berkeley Wireless Research Center at the University of California at Berkeley, where he is working on the development of the human intranet. His research focused on personal exposure assessment to radio-frequency electromagnetic fields and numerical dosimetry. He was a recipient of the Joseph James Morrissey Memorial Award issued by the Bioelectromagnetics Society (BEMS) and the European Bio Electromagnetics Association (EBEA) in June 2013. He received the FWO [PEGASUS] Marie Skłodowska-Curie Fellowship in 2017, funded by the European Union's Horizon 2020 Research and Innovation Programme under the Marie Skłodowska-Curie Grant Agreement 665501, with the research Foundation Flanders (FWO). He also received the Honorary Fellowship by the Belgian American Education Foundation (BAEF) in 2017, and the 2015 International Union of Radio Science (Union Radio-Scientifique Internationale or URSI) Young Scientist Award.



Günter Vermeeren (Member, IEEE) received the M.Sc. degree in industrial engineering from the KAHO Sint-Lieven, Ghent, Belgium, in July 1998, and the M.Sc. degree in electrical engineering and the Ph.D. degree in electro-technical engineering from Ghent University, Belgium, in July 2001 and August 2013, respectively. From September 2001 to September 2002, he joined the Research and Development Department of network integrator Telindus, Leuven, Belgium. Since September 2002, he has been a Research Engineer with the WiCa group

of Prof. L. Martens. His research interests include the numerical modeling as well as measurements of electromagnetic radiation in the domain of radio-frequency dosimetry, electromagnetic exposure, on-body propagation, interference in communication networks, and medical imaging systems, such as, hybrid MRI systems.







Luc Martens (Member, IEEE) received the M.Sc. degree in electrical engineering from Ghent University, Ghent, Belgium, in 1986, and the Ph.D. degree in 1990. From September 1986 to December 1990, he was a Research Assistant with the Department of Information Technology (INTEC), Ghent University. During this period, his scientific work focused on the physical aspects of hyperthermic cancer therapy. His research work dealt with electromagnetic and thermal modeling, and the development of measurement systems for that application. Since 1991,

he manages the WAVES research group (at the time; the Wireless and Cable Research Group), INTEC. This group is part of IMEC Institute since 2004. Since April 1993, he has been a Professor with Ghent University. He has authored/coauthored more than 300 publications in the domain of electromagnetic channel predictions, dosimetry, exposure systems and health, and wireless communications. His research interests include modeling and measurement of electromagnetic channels, electromagnetic exposure, e.g., around telecommunication networks and systems, such as cellular base station antennas and energy consumption of wireless networks.



Guy Torfs (Senior Member, IEEE) received the Engineering degree in applied electronics and the Ph.D. degree in applied sciences and electronics from Ghent University, Ghent, Belgium, in 2007 and 2012, respectively. Since 2011, he has been with IMEC, associated with Ghent University, where he became an Assistant Professor in 2015. His research interests include high-speed mixed-signal designs for wireless baseband and fiber-optic and backplane communication systems, including digital signal processing and calibration, analog equalization

circuits, and clock and data recovery systems. In 2014, as part of the Bi-PON and Cascaded Bi-PON Team, he was a recipient of the Greentouch 1000x Award. He was a co-recipient of the 2015 DesignCon Best Paper Award in the High-Speed Signal Design category. He serves as an Associate Editor of the IEEE TRANSACTIONS ON CIRCUITS AND SYSTEMS- II and IEICE Elex.



Wout Joseph (Senior Member, IEEE) was born in Ostend, Belgium, in October 1977. He received the M.Sc. degree in electrical engineering from Ghent University, Ghent, Belgium, in 2000, and the Ph.D. degree in electrical engineering, Ghent University, Ghent, in 2005. From September 2000 to March 2005, he was a Research Assistant with the Department of Information Technology (INTEC), Ghent University. During this period, his scientific work focused on electromagnetic exposure assessment. Since April 2005, he has been a Post-Doctoral

Researcher with imec-UGent/INTEC. From October 2007 to October 2013, he was a Post-Doctoral Fellow of FWO-V (Research Foundation-Flanders). Since October 2009, he has been a Professor in the domain of experimental characterization of wireless communication systems. His research interests include measuring and modeling electromagnetic fields around base stations for mobile communications, health effects of exposure to electromagnetic radiation, electromagnetic exposure assessment, propagation for wireless communication systems, and antennas and calibration. Furthermore, he specializes in wireless performance analysis and quality of experience.